

The $X(3872)$ at CDF II

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Last year's $X(3872)$ discovery was confirmed with the CDF II detector in $\bar{p}p$ collisions. We measure its mass to be $3871.3 \pm 0.7 \pm 0.4$ MeV/ c^2 . The source of X -mesons in the large CDF sample is resolved by studying their vertex displacement. We find $16.1 \pm 4.9 \pm 2.0\%$ of our X -sample comes from decays of b -hadrons, and the remainder from prompt sources: either direct production or by decay of (unknown) short-lived particles. The mix of production sources is similar to that observed for the $\psi(2S)$ charmonium state.

Keywords: $X(3872)$; Charmonium.

At last year's Lepton-Photon Symposium Belle announced discovery of a charmonium-like state,^{1,2} $X(3872)$, in $B^+ \rightarrow K^+ J/\psi \pi^+ \pi^-$. CDF quickly confirmed $X \rightarrow J/\psi \pi^+ \pi^-$.³ A natural interpretation of the X is the 3D_2 of $c\bar{c}$, but this is contrary to expectations. The 3D_2 is thought to be significantly lighter (~ 3830 MeV/ c^2); and Belle failed to detect decays to $\chi_{c1}\gamma$, which should be prominent for 3D_2 . More circumstantial is the expectation of a relatively flat dipion mass ($M_{\pi\pi}$) distribution for D -states,⁴ whereas Belle found high masses preferred—possibly consistent with the (isospin violating) decay to $J/\psi \rho^0$. These difficulties, coupled with the proximity of the $X(3872)$ to the $D^0 \bar{D}^{*0}$ -threshold, prompted speculation that the X may be a $D^0 \bar{D}^{*0}$ “molecule.” Whether this is the case, or the X is “only” a $c\bar{c}$ state in conflict with current theoretical models, the X is an interesting object of study.⁵

CDF II⁶ is a general purpose detector at Fermilab's $\bar{p}p$ collider. We use 220 pb^{-1} of $\mu^+ \mu^-$ triggers, yielding a clean J/ψ sample. Aside from technical cuts, kinematic and spatial cuts suppress large backgrounds from J/ψ 's plus random tracks. The main cuts are: a maximum number of $J/\psi \pi \pi$ candidates/event, $p_T(J/\psi) > 4$ GeV/ c , $p_T(\pi) > 400$ MeV/ c , and $\Delta R \equiv \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2} < 0.7$ for each pion, where $\Delta\phi$ ($\Delta\eta$) is the azimuthal (pseudorapidity) difference of the pion with respect to the $J/\psi \pi \pi$. With these cuts a significant X -signal is revealed.³ Here, however, we show in Fig. 1 the results split up into $M_{\pi\pi} < 500$ and > 500 MeV/ c^2 subsamples. No X -signal is apparent for low $M_{\pi\pi}$, supporting Belle's observation of high-mass decays.

Using the high- $M_{\pi\pi}$ sample, the X -mass is $3871.3 \pm 0.7 \pm 0.4$ MeV/ c^2 . Also shown in Fig. 1 are masses from other experiments, and the average compared to the $D^0 \bar{D}^{*0}$ threshold. The near equality helps fuel molecular- $D^0 \bar{D}^{*0}$ speculations.

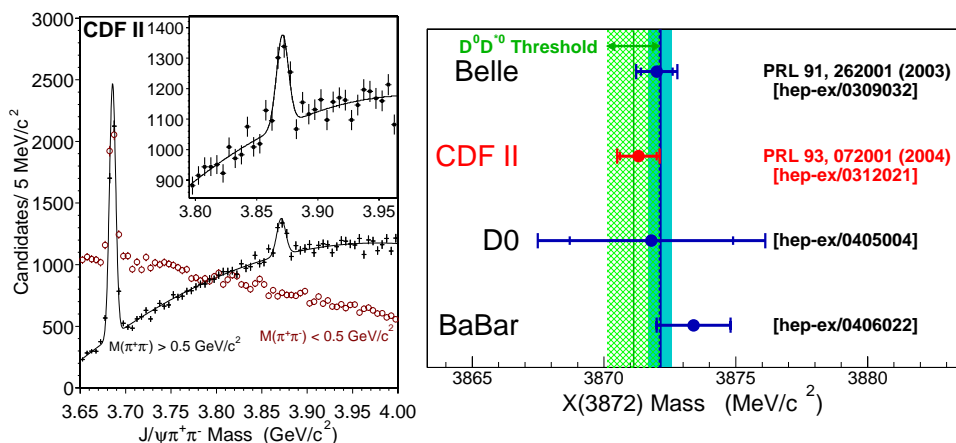


Fig. 1. **LEFT:** The $J/\psi\pi^+\pi^-$ mass distribution for $M_{\pi\pi} < 500$ and > 500 MeV/c^2 subsamples. **RIGHT:** Summary of X -mass measurements compared to the $D^0\bar{D}^{*0}$ threshold.

From Belle’s observation, B -mesons are a significant source of X ’s. This raises some questions: Is the CDF sample only from b -hadrons? If not, is direct X production different from charmonium? The technique of separating b -decay feeddown from prompt sources is well established.⁷ Since X -decay is not weak, it is too rapid to leave a displaced vertex. If, however, it is produced by a (boosted) b -decay, the X will be displaced due to the b -lifetime. We measure the transverse X -displacement, L_{xy} , and convert it to an “uncorrected” proper-time: $ct = M \cdot L_{xy}/p_T$. This is not the true proper-time of the b -decay because M and p_T are only for the $J/\psi\pi^+\pi^-$.

We use the same X -sample as above, but now impose additional cuts related to the Si -vertex tracker, mainly to demand $\sigma(L_{xy}) < 125 \mu\text{m}$ and have good beamline information. The sample is reduced by $\sim 15\%$. An unbinned likelihood fit is performed simultaneously over the ct and mass of the candidates. The signal is modeled by a Gaussian in mass; and for the ct -distribution, a resolution smeared exponential for the long-lived component and by the resolution function for the prompt. The background model uses a quadratic polynomial for mass, and resolution function for the prompt and *three* resolution smeared exponentials—one for the negative- ct tail and two for the positive. The resolution function consists of two Gaussians.

The fit for $\psi(2S)$ is shown in Fig. 2, where $28.3 \pm 1.0 \pm 0.7\%$ of signal is displaced, similar to Run I results.⁷ For the $X(3872)$, with $M_{\pi\pi} > 500$ MeV/c^2 , the fraction is $16.1 \pm 4.9 \pm 2.0\%$ (Fig. 3)—a bit more than 2σ from the $\psi(2S)$. These fractions agree with those obtained by simple sideband subtraction. They are, however, uncorrected for efficiency, and must be considered sample specific.⁷ The *absence* of a b -component is excluded at 3σ based on Monte Carlo “experiments.” Thus our X -sample is mainly prompt—presumably direct production—with a modest b -contribution.

It has been argued that all conventional $c\bar{c}$ assignments for the $X(3872)$ are

Fig. 2. Projection of $\psi(2S)$ likelihood fit onto the uncorrected proper-time distribution for the full PDF, and its breakdown into signal (shaded) and background (hatched) classes. Signal and background are further separated into prompt and long-lived components. The projections are for candidates within $\pm 2.5\sigma$ of the $\psi(2S)$ mass in order to be reflective of its signal-to-background ratio. The fit actually spans the mass range 3640-3740 MeV/c^2 .

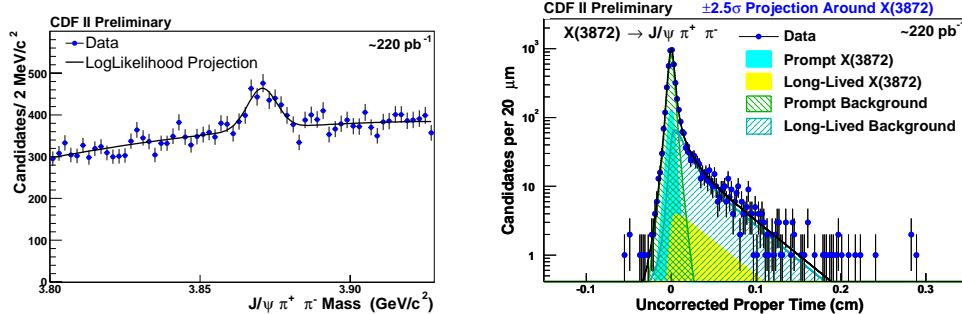


Fig. 3. Projections of X -likelihood fit in mass (left), and uncorrected proper-time (right) as Fig. 2.

problematic.⁸ However, production of the X appears, so far, quite similar to that of the $\psi(2S)$ in CDF. If it is indeed a “molecule,” there seems to be no dramatic penalty for producing such a fragile state in $\bar{p}p$ collisions. Although, more incisive comparisons require specific theoretical models for the production of exotic states. A recent analysis of X -production as a 1^{++} state⁹ may benefit from our results.

Studies of this mysterious state are continuing in CDF.

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